### nature communications



**Article** 

https://doi.org/10.1038/s41467-022-32622-4

# High-surface-area corundum nanoparticles by resistive hotspot-induced phase transformation

Received: 9 July 2021

Accepted: 9 August 2022

Published online: 26 August 2022

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High-surface-area α-Al<sub>2</sub>O<sub>3</sub> nanoparticles are used in high-strength ceramics and stable catalyst supports. The production of α-Al<sub>2</sub>O<sub>3</sub> by phase transformation from γ-Al<sub>2</sub>O<sub>3</sub> is hampered by a high activation energy barrier, which usually requires extended high-temperature annealing (~1500 K, > 10 h) and suffers from aggregation. Here, we report the synthesis of dehydrated α-Al<sub>2</sub>O<sub>3</sub> nanoparticles (phase purity ~100%, particle size ~23 nm, surface area ~65 m<sup>2</sup> g<sup>-1</sup>) by a pulsed direct current Joule heating of γ-Al<sub>2</sub>O<sub>3</sub>. The phase transformation is completed at a reduced bulk temperature and duration (~573 K, < 1 s) via an intermediate δ'-Al<sub>2</sub>O<sub>3</sub> phase. Numerical simulations reveal the resistive hotspot-induced local heating in the pulsed current process enables the rapid transformation. Theoretical calculations show the topotactic transition (from γ- to δ'- to α-Al<sub>2</sub>O<sub>3</sub>) is driven by their surface energy differences. The α-Al<sub>2</sub>O<sub>3</sub> nanoparticles are sintered to nanograined ceramics with hardness superior to commercial alumina and approaching that of sapphire.

High-surface-area corundum nanoparticles ( $\alpha$ -Al<sub>2</sub>O<sub>3</sub> NPs) have widespread applications. For examples, corundum is widely used in ceramics for biomedical implants<sup>1,2</sup> and high-speed cutting tools<sup>3</sup>.  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> NPs precursors provide access to nanometer-grained alumina ceramics with significantly improved fracture toughness<sup>4</sup>, wear resistance<sup>5</sup>, and high density under reduced sintering temperature<sup>6</sup>. Even though  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> NPs are primarily used as catalyst supports due to their high surface areas<sup>7</sup>, the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> with high surface area could be used as a catalytic support in auto-exhaust Pt-Mo-Co catalytic converters<sup>8</sup> and enhance Ru catalyst activity for ammonia synthesis<sup>9</sup>. The high mechanical stability of  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> enables a low sintering behavior, which is vital for its use in reforming reactions to obtain synthesis gas under harsh conditions<sup>10,11</sup>.

Much effort has been made toward improving the synthesis of  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>, yet very few of the processes afford high-surface-area NPs due to various intrinsic thermodynamic limits<sup>6,12,13</sup>. First, even though corundum is the thermodynamically stable phase of coarsely crystalized aluminum oxide (Al<sub>2</sub>O<sub>3</sub>), the synthesis of nanocrystalline Al<sub>2</sub>O<sub>3</sub> usually leads to  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> because of its lower surface energy based on previous experimental observation and theoretical calculation <sup>12,14,15</sup>. The second reason is the high activation energy barrier of -485 kJ mol $^{-1}$  for the phase transformation from the cubic close-packed structure of the  $\gamma$ -phase to the hexagonal close-packed structure of the  $\alpha$ -phase that involves intensive bond breaking and remaking <sup>16</sup>. Thirdly, the density of  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> (3.99 g cm $^{-3}$ )<sup>17</sup> is higher than that of the transition alumina phases (3.6–3.67 g cm $^{-3}$  for  $\gamma$ -,  $\eta$ -, and  $\delta$ -Al<sub>2</sub>O<sub>3</sub>)<sup>17</sup>, thus sufficient energy

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or high pressure<sup>18</sup> is needed for the density uphill process from transition alumina to  $\alpha\text{-}Al_2O_3$ . As a result, the thermal processes usually require temperatures >1470 K with prolonged annealing times of 10 to 20 h to facilitate the phase transformation  $^{16,17}$ , which could also result in detrimental aggregation and sintering of alumina phases. The high-energy input and extended high-temperature annealing usually lead to surface area <10 m² g⁻¹ because of the substantial mass transfer¹³. Moreover, the polymorphism of  $Al_2O_3$  during the phase transformation further increases the complexity and could lead to the mixed transition alumina with undesired  $\delta\text{-}$  and  $\theta\text{-}Al_2O_3^{16,19,20}$ .

The phase transformation and grain growth of alumina usually coexist in any thermal process. Previous kinetics studies show that the activation energies of the phase transformation from  $\gamma$ - to  $\alpha\text{-Al}_2O_3$  are around  $200\text{--}500\,\text{kJ}\,\text{mol}^{-1}$  depending on the sample conditions  $^{16,21,22}$ ; in contrast, the activation energies for the grain growth are around  $500\text{--}900\,\text{kJ}\,\text{mol}^{-1}$  depending on the grain-boundary orientations  $^{23\text{--}26}$ . This implies that the kinetics of the phase transformation is possibly faster than the grain growth. Hence, based on these prior results, it is reasonable to presume that a rapid and lower-temperature thermal process would reduce the grain coarsening and maintain high surface area during the phase transformation process.

To meet these goals, here, we show a Joule heating process based on pulsed direct current (PDC) to complete the phase transformation from y- to α-Al<sub>2</sub>O<sub>3</sub> at a significantly reduced average bulk temperature and reaction duration (~573 K, <1s). The rapid transformation is enabled by the resistive hotspot-induced local heating in the PDC process when an appropriate volume fraction ratio of γ-Al<sub>2</sub>O<sub>3</sub> precursors and carbon black conductive additives are used. The pulsed and local heating mitigates the aggregation, leading to the synthesis of α-Al<sub>2</sub>O<sub>3</sub> NPs with an average particle size of ~23 nm and a surface area of ~65 m<sup>2</sup> g<sup>-1</sup>. Ab initio calculations reveal that the topotactic phase transformation process (from y- to  $\delta'$ - to  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>) is driven by the surface energy difference of the three phases. The calculations suggest that a particle size of ~17 nm is the thermodynamic limit for the synthesis of anhydrous  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> NPs with the  $\delta$ '-Al<sub>2</sub>O<sub>3</sub> as the intermediate phase by a thermal process, matching well with the experimental values. Also, based on the Joule heating technique, we develop an alternating current sintering (ACS) process and show the ultrafast and pressureless sintering of the α-Al<sub>2</sub>O<sub>3</sub> NPs into alumina ceramics with nanoscale grain size. The ceramics from these α-Al<sub>2</sub>O<sub>3</sub> NPs by twostep pressureless sintering process demonstrate hardness of -15 GPa, superior to commercial standard alumina and comparable to singlecrystal sapphire.

#### Results

## Phase transformation synthesis of corundum nanoparticles by PDC

Since the y-Al<sub>2</sub>O<sub>3</sub> NPs precursors are electrically insulative, commercial carbon black (CB) was used as the conductive additive. In a typical experiment, the mixture of γ-Al<sub>2</sub>O<sub>3</sub> NPs and CB were compressed inside a quartz tube between two graphite electrodes (Fig. 1a, Supplementary Fig. 1a and Supplementary Table 1). The CB also works as separators to avoid the aggregation of Al<sub>2</sub>O<sub>3</sub> NPs during heating. The carbon black is composed of ultrafine amorphous carbon nanospheres and has a surface area of ~1600 m<sup>2</sup> g<sup>-1</sup>, which permits the intimate mixing with γ-Al<sub>2</sub>O<sub>3</sub> precursors, as confirmed by the energy dispersive X-ray spectroscopy (EDS) maps (Supplementary Fig. 2). The resistance was controlled by the compressive force on the two electrodes (Supplementary Table 1). The electrodes were connected to a capacitor bank with total capacitance of C = 0.624 F and charging voltage up to  $V_0$  = 500 V. The discharge circuit is a series resistor-inductor-capacitor circuit with the characteristic time of  $\tau = 0.1$  ms, which permits the PDC with frequency of  $f = 1000 \, \text{Hz}$  (Supplementary Fig. 1b). Joule heating affects the entire electric conductor; for a homogeneous conductor, the current density is uniform so the Ohmic dissipation enables the homogeneous temperature distribution throughout the sample<sup>27</sup>. However, when an electrical field is applied to an inhomogeneous medium, as in the composite of conductive CB and insulative  $Al_2O_3$ , the current and powder densities have strong spatial variation<sup>28</sup>. The power dissipation is substantially larger at some regions than the neighbor; these regions are termed resistive hotspots (Fig. 1a). Even though the average bulk temperature is low, the hotspots permit local heating and trigger the transformation that happens at a much higher temperature.

By using this effect, we realized the phase transformation from y- $Al_2O_3$  to  $\alpha$ - $Al_2O_3$  accompanied by the intermediate phase of  $\delta'$ - $Al_2O_3$  at an average bulk temperature of ~573 K in <1 s (Fig. 1b, bottom). We compared our process to the representative phase transformation methods reported in the literatures (Fig. 1b)<sup>13,16,20</sup>. The liquid-feed flame spray pyrolysis produces α-Al<sub>2</sub>O<sub>3</sub> at temperatures near 1873 K (Fig. 1b, top); however, the kinetically controlled process may render it difficult to access the pure phase  $(80-85\% \text{ purity of } \alpha\text{-phase})^{20}$ . This phase purity is not an issue for ceramics since all other phases would transform to α-phase during the sintering process; nevertheless, the pure phase would be important for other applications such as in catalyst supports. Traditional heating methods that supply heat through the sample boundary, such as furnace annealing, require an extended period to permit uniform heating; hence 1473 K and 10 to 20 h are necessary to complete the phase conversion (Fig. 1b, middle)<sup>16</sup>. Other room-temperature nonequilibrium processes, such as high-energy ball milling, have been reported to form  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> (Fig. 1b, middle)<sup>13</sup> that shows hydrothermal stability<sup>29</sup>. Nevertheless, the Al<sub>2</sub>O<sub>3</sub> can aggregate, which leads to loss of surface area during the extended time and highenergy collisions<sup>30,31</sup>.

We investigated the detailed phase transformation process of y-Al<sub>2</sub>O<sub>3</sub> by the PDC approach (Fig. 1c, d, e). Commercial y-Al<sub>2</sub>O<sub>3</sub> NPs with particle size of 5-10 nm and surface area of ~156 m<sup>2</sup> g<sup>-1</sup> were used as the precursors (Supplementary Figs. 3-5). X-ray diffraction (XRD) characterization and Rietveld refinement show that the starting materials are composed of ~91 wt% y-Al<sub>2</sub>O<sub>3</sub> with crystalline size of ~4 nm, and ~9 wt% y-Al(OH)<sub>3</sub> (Supplementary Fig. 6a). The y-Al(OH)<sub>3</sub> could be easily decomposed to γ-Al<sub>2</sub>O<sub>3</sub> by mild calcination (Supplementary Fig. 6b). The mass ratio of γ-Al<sub>2</sub>O<sub>3</sub> NPs and CB is 4 to 1, which gives a sample resistance of ~8 Ω (Supplementary Table 1). A discharging voltage of 60 V was applied with different discharging times controlled by a relay. The XRD patterns of the products with different PDC onstate time are shown in Fig. 1c. As the discharging time increased, the y-Al(OH)<sub>3</sub> first disappeared at 0.3 s; then, the γ-Al<sub>2</sub>O<sub>3</sub> was transferred to the mixed  $\delta'$ - and  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> at 0.4 to 0.5 s; last, the intermediate  $\delta'$ -Al<sub>2</sub>O<sub>3</sub> was fully converted to α-Al<sub>2</sub>O<sub>3</sub> after 0.8 s of discharge (Fig. 1e). The orthorhombic δ'-Al<sub>2</sub>O<sub>3</sub> is observed as the single intermediate phase (Fig. 1d), which is distinct from other thermal processes where  $\delta$ - and  $\theta$ - $Al_2O_3$  usually appear before the final  $\alpha$ - $Al_2O_3$  phase (Fig. 1b)<sup>16,17,32</sup>.

Unlike our previous report<sup>33</sup> on the synthesis of graphene by the high-voltage flash Joule heating at a high temperature of ~3000 K, the 60 V PDC does not provide enough energy to graphitize the CB (Supplementary Fig. 7). As a result, the residual CB (~20 wt%) could be easily removed by heating in air, according to thermal gravimetric analysis (TGA) (Supplementary Fig. 8a). Here, the as-synthesized mixture of α-Al<sub>2</sub>O<sub>3</sub> NPs and CB was calcined in air at 700 °C for 1 h to purify the product. The X-ray photoemission spectrum (XPS) of the α-Al<sub>2</sub>O<sub>3</sub> product after calcination showed very minor carbon signal, which could be caused by the carbon adsorption in air (Supplementary Fig. 8b). Raman spectra are sensitive to even a monolayer of carbon<sup>34</sup>; intriguingly, no characteristic Raman bands of carbon were detected after calcination at 700 °C (Fig. 1f), demonstrating the efficient removal of carbon. As a control, we show that the calcination process itself does not trigger the phase transformation and has negligible effect on the coarsening or aggregation of the y-Al<sub>2</sub>O<sub>3</sub> phase (Supplementary Fig. 9).

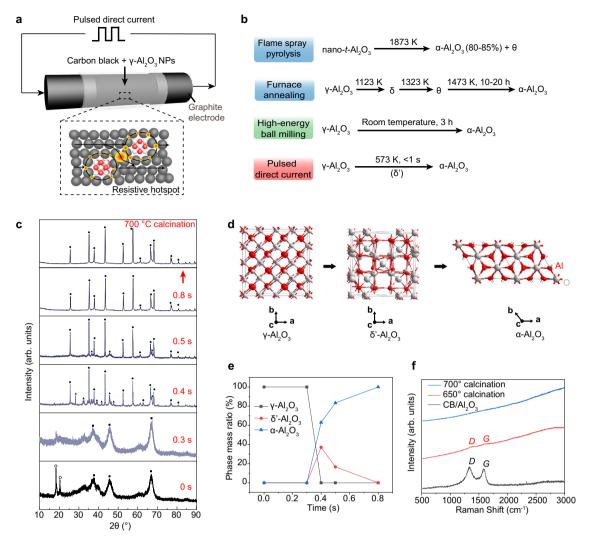


Fig. 1 | Ultrafast phase transformation of alumina by pulsed direct current (PDC) Joule heating. a Schematics of the PDC apparatus, and the resistive hotspots around and at the gaps of the insulative  $\gamma\text{-}Al_2O_3$  NPs. The black arrows depict the electric current lines. b Representative methods for the phase transformation from  $\gamma$ - to  $\alpha\text{-}Al_2O_3$ : flame spray pyrolysis, ref. 20 furnace annealing, ref. 16 high-energy ball milling, ref. 13 PDC, this work. c X-ray diffraction (XRD) patterns of  $\gamma\text{-}Al_2O_3$  NPs after different PDC durations and the  $\alpha\text{-}Al_2O_3$  NPs product after calcination. The marks:  $\gamma\text{-}Al_2O_3$  (square),  $\delta\text{-}Al_2O_3$  (triangle),  $\alpha\text{-}Al_2O_3$  (dot), and  $\gamma\text{-}Al(OH)_3$  (circle). The precursor is  $\gamma\text{-}Al_2O_3$  with -9 wt%  $\gamma\text{-}Al(OH)_3$  phase (gibbsite, crystal system:

monoclinic; space group: P21/n; PDF No. 07-0324). The 0.8 s sample was calcined at 700 °C for 1 h. **d** Crystal structures of alumina phases:  $\gamma\text{-}Al_2O_3$  (crystal system: cubic; space group: Fd-3m; PDF No. 10-0425),  $\delta\text{'-}Al_2O_3$  (crystal system: orthorhombic; space group: P222; PDF No. 46–1215), and  $\alpha\text{-}Al_2O_3$  (crystal system: trigonal; space group: R-3c; PDF No. 46–1212). For  $\gamma\text{-}Al_2O_3$ , all the Al sites are depicted to show the crystal structure, while in the actual structure, not all the sites are 100% occupied. **e** Phase mass ratio of alumina polymorphs varied with PDC duration. **f** Raman spectra of as-synthesized  $\alpha\text{-}Al_2O_3$ /carbon black mixture and the purified  $\alpha\text{-}Al_2O_3$  NPs by calcination at different temperatures.

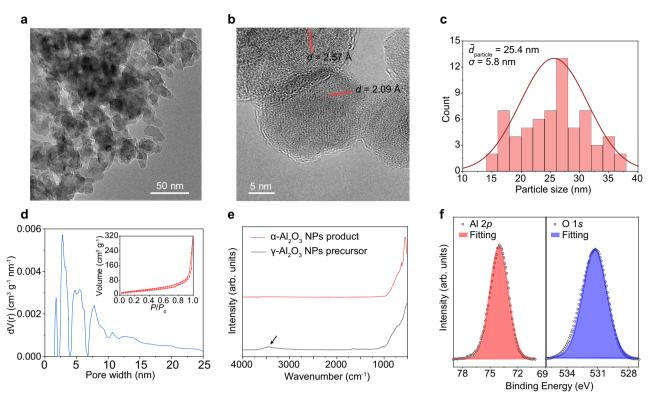
#### Characterization of the corundum nanoparticles

The α-Al<sub>2</sub>O<sub>3</sub> NPs derived by PDC followed by mild calcination were further characterized in detail. Bright-field transmission electron microscopy (BF-TEM) images showed the well-dispersed particles (Fig. 2a). High-resolution TEM (HRTEM) showed the high degree of crystallinity of the α-Al<sub>2</sub>O<sub>3</sub> NPs (Fig. 2b). The interplanar spacing values of ~2.57 Å and ~2.09 Å correspond to the  $d(10\overline{1}4)$  and  $d(11\overline{2}3)$ of α-Al<sub>2</sub>O<sub>3</sub>, respectively. The single set of lattice fringes in the NP and the nanobeam diffraction (NBD) show that the as-synthesized α-Al<sub>2</sub>O<sub>3</sub> NPs are single-crystal particles (Supplementary Fig. 10). We observed some α-Al<sub>2</sub>O<sub>3</sub> NPs with surface roughness features at a few nm (Supplementary Fig. 10), which are similar to the particle size of the y-Al<sub>2</sub>O<sub>3</sub> precursors (Supplementary Fig. 4). This demonstrates that the rapid PDC process triggers the phase transformation while no significant aggregation of the NPs occurs. The TEM images show that the particle size ranged from 14 to 36 nm, with an average particle size of 25.4 nm and standard derivation ( $\sigma$ ) of 5.8 nm (Fig. 2c).

Brunauer–Emmett–Teller (BET) measurement showed that the surface area of the  $\alpha$ -Al $_2$ O $_3$  NPs is -65 m $^2$  g $^{-1}$  (Fig. 2d, inset). The average particle size (D) is estimated to be -23 nm by Eq. (1),

$$D = 6/(\rho S) \tag{1}$$

where  $\rho$  is the density of  $\alpha$ -Al $_2$ O $_3$  (3.96 g cm $^{-3}$ ) and S is the specific surface area $^{35}$ . The pore size determined from the N $_2$  adsorption-desorption isotherm using the density functional theory (DFT) model $^{36}$  indicates the distribution with high probability at 3 to 10 nm (Fig. 2d and Supplementary Fig. 11). The observed surface area was mainly attributed to the nanoscale grain size or intraparticle pores, and minor interparticle voids and surface roughness features of the NPs (Supplementary Fig. 10). The crystalline size of the  $\alpha$ -Al $_2$ O $_3$  NPs was estimated to be -22 nm based on the Halder-Wagner method (Supplementary Note 1 and Supplementary Fig. 12). The crystalline size (-22 nm) agrees well with the particle size measured from TEM statistics (-25 nm) and BET estimation (-23 nm), demonstrating the



**Fig. 2** | **Characterization of the α-Al<sub>2</sub>O<sub>3</sub> NPs. a** Bright-field transmission electron microscopy (BF-TEM) image of the α-Al<sub>2</sub>O<sub>3</sub> NPs. **b** High-resolution TEM (HRTEM) image of the α-Al<sub>2</sub>O<sub>3</sub> NPs. The *d* spacing of 2.57 Å and 2.09 Å correspond to the  $d(10\overline{1}4)$  and  $d(11\overline{2}3)$  of α-Al<sub>2</sub>O<sub>3</sub>. **c** Histogram and distribution of the α-Al<sub>2</sub>O<sub>3</sub> NPs particle size determined by TEM. **d** Pore width distribution determined by the

application of density functional theory (DFT) model to the  $N_2$  isotherm. Inset,  $N_2$  adsorption–desorption isotherm of as-synthesized  $\alpha\text{-}Al_2O_3$  NPs at 77 K.  $\boldsymbol{e}$  Fourier-transform infrared (FT-IR) spectra of the  $\gamma\text{-}Al_2O_3$  NPs precursor and the  $\alpha\text{-}Al_2O_3$  NPs product. The black arrow points to the hydroxyl group absorbance.  $\boldsymbol{f}$  X-ray photoelectron spectroscopy (XPS) fine spectra of Al and O of the  $\alpha\text{-}Al_2O_3$  NPs.

single-crystal feature of the NPs. The dynamic light scattering (DLS) measurement shows that the as-synthesized  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> NPs are well-dispersible (Supplementary Fig. 13).

Unlike the starting  $\gamma\text{-}Al_2O_3$  NPs that have hydrated surface states, the as-synthesized  $\alpha\text{-}Al_2O_3$  NPs surfaces are highly dehydrated because of the thermal process (Fig. 2e). The XPS fine spectra showed the dominate  $O^2$  peak at a binding energy of ~531.2 eV and single  $Al^{3+}$  peak at a binding energy of ~74.0 eV from the  $\alpha\text{-}Al_2O_3$  NPs (Fig. 2f). This demonstrated that the ultrafast PDC process does not result in obvious oxygen deficiencies or the carbothermic reduction of  $Al_2O_3$  even with the existence of CB, presumably due to the high reduction potential of  $Al^{3+}$ . No other peak was detected in the XPS full spectrum (Supplementary Fig. 8b), indicating the high-purity synthesis ability of the electric thermal process. This makes it superior to the solvent-based methods including ball milling  $^{13}$  or co-precipitation  $^6$ , which inevitably suffer from lengthy purification processes and chemical contaminants.

#### Resistive hotspot effect revealed by current density simulation

As discussed before, the composition of an inhomogeneous media is critical for local power dissipation during the PDC process. To quantitatively investigate the effect of the composition on the phase transformation, a series of precursors with different mass ratio of  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> and CB were treated by PDC under the same voltage and time (Fig. 3a and Supplementary Table 1). According to the densities of  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> and CB, the volume fractions (f) of  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> are obtained (Supplementary Note 2 and Supplementary Table 2), and the phase mass ratios varied with  $f(\gamma$ -Al<sub>2</sub>O<sub>3</sub>) after the PDC process are calculated (Fig. 3b). The phase transformation degree is increased as the  $f(\gamma$ -Al<sub>2</sub>O<sub>3</sub>) increased from 0.41 to 0.73; the phase-pure  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> is obtained at  $f(\gamma$ -Al<sub>2</sub>O<sub>3</sub>) -0.73. Further increase in the  $f(\gamma$ -Al<sub>2</sub>O<sub>3</sub>) to  $\geq$  0.78 leads to no phase transformation.

To explain the  $f(\gamma - Al_2O_3)$ -dependent phase transformation, the electrical conductivity and temperature were measured. The conductivities are calculated based on the measured resistance (R) and the feature size of the samples (Supplementary Table 1 and Fig. 3c). The conductivity is inversely proportional to  $f(\gamma - Al_2O_3)$ , which is reasonable since  $\gamma - Al_2O_3$  is electrically insulative. The real-time temperature was measured using an infrared (IR) thermometer (Supplementary Fig. 14). The average bulk temperature is decreased with the increase of  $f(\gamma - Al_2O_3)$  (Fig. 3c). This could be explained by the power (P) equation of Joule heating by Eq. (2),

$$P = \frac{V^2}{R} = V^2 \sigma \tag{2}$$

where V is the voltage, and  $\sigma$  is the conductivity of the sample. Since the start voltages were fixed to  $V_0$  = 60 V, the power was proportional to the conductivity of the sample. Intriguingly, the phase-pure  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> NPs were obtained at a low average bulk temperature of ~573 K with  $f(\gamma$ -Al<sub>2</sub>O<sub>3</sub>) ~0.73 (Fig. 3c).

Such a low temperature is not supposed to trigger the phase transformation from  $\gamma$ - to  $\alpha\text{-Al}_2O_3$  with a high activation energy of ~485 kJ mol $^1$  (ref. 16). The phase transformation temperatures from transition alumina to  $\alpha\text{-Al}_2O_3$  in other thermal processes are substantially higher (Fig. 1b), e.g., flame spray pyrolysis at 1873 K (ref. 20), furnace annealing at 1473 K (ref. 16), and the annealing process even with  $\alpha\text{-Fe}_2O_3$  seeds at 973 K (ref. 6). Moreover, the higher phase transformation degree at a lower temperature is counterintuitive (Fig. 3c). To explain the intriguing phenomenon, we conduct a numerical simulation based on the finite element method (FEM) on the current density distribution of the  $\gamma\text{-Al}_2O_3/\text{CB}$  composite during PDC process (see details in Supplementary Note 2, Supplementary

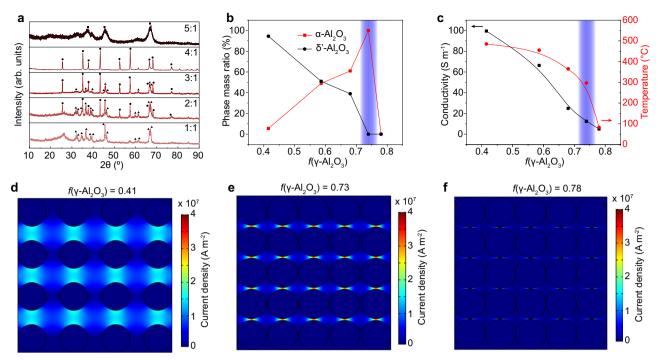


Fig. 3 | Resistive hotspot effect in pulsed direct current (PDC) process. a X-ray diffraction (XRD) patterns of  $\gamma\text{-}Al_2O_3/CB$  with different mass ratios after the same PDC process. The marks are  $\gamma\text{-}Al_2O_3$  (square),  $\delta\text{-}Al_2O_3$  (triangle), and  $\alpha\text{-}Al_2O_3$  (dot). The numbers are the mass ratio of  $\gamma\text{-}Al_2O_3$  to carbon black (CB). b Phase mass ratios of the product after PDC process varied with volume fraction of  $\gamma\text{-}Al_2O_3$ ,  $f(\gamma\text{-}Al_2O_3)$ . The blue region denotes the pure  $\alpha\text{-}phase$  at f=0.73. c Conductivities and

temperatures varied with  $f(\gamma \cdot Al_2O_3)$ . The blue region denotes the pure  $\alpha$ -phase at f=0.73.  $\mathbf{d}-\mathbf{f}$  Simulated current density maps of the sample during PDC with different  $\gamma \cdot Al_2O_3$  volume fractions of  $\mathbf{d}f=0.41$ ,  $\mathbf{e}f=0.73$ , and  $\mathbf{f}f=0.78$ . The separated balls are  $\gamma \cdot Al_2O_3$  NPs and the continuous phase is CB. The color bars show the current density values.

Figs. 15–19, and Supplementary Tables 3–5). As shown in Fig. 3d–f, the current density is inhomogeneous in the composite of  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> and CB; the current densities at the regions of vertical gaps between  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> NPs are larger than the bulk regions. The gaps become narrower as the  $f(\gamma$ -Al<sub>2</sub>O<sub>3</sub>) increased, leading to substantially large current densities in those regions. Considering that the resistivity (R) of the conductive CB phase is constant, the heat (Q) per volume produced by PDC is proportional to the square of the current density (f) by Eq. (3),

$$Q \propto j^2 R \tag{3}$$

The large thermal dissipation in the regions with high current densities leads to the hotspots near  $\gamma\text{-}Al_2O_3$  NPs with much higher temperature than the bulk regions, which triggers the phase transformation.

Experimentally,  $f(\gamma-Al_2O_3)$  of ~0.73 is the optimized volume fraction for the rapid and thorough phase transformation. Since the phase transformation temperature from y- to  $\alpha$ -phase using a furnace annealing is reported to be ~1473 K (ref. 16), we here defined the region as the hotspot zone with temperature  $T_{\text{hotspot}} \ge 1473 \text{ K}$ . Under such a definition, the hotspot zone is depicted (Supplementary Fig. 18), and it is estimated that ~30% of the particle surface area is heated to above the phase transformation temperature (see details in Supplementary Note 2, Estimation of the temperature and zone size of the hotspot). The quantitative analysis of the current densities suggests a decreased bulk temperature but an increased hotspot temperature as the  $f(\gamma-Al_2O_3)$  increases (Supplementary Fig. 19), which agrees well with the average bulk temperature measurement (Fig. 3c). Thus, the resistive hotspot-induced local heating during the PDC process well explains the observed phase transformation at a low bulk average temperature.

#### Topotactic transition pathway revealed by ab initio calculations

To provide deeper insight into the topotactic transition pathway, we conducted thermodynamic analysis of the three Al<sub>2</sub>O<sub>3</sub> phases based on DFT (see details in Methods). Both the bulk energies and surface energies of the three Al<sub>2</sub>O<sub>3</sub> phases were calculated (Fig. 4a and Supplementary Table 6). The bulk energy of α-Al<sub>2</sub>O<sub>3</sub> is the lowest, followed by that of  $\delta'$ -Al<sub>2</sub>O<sub>3</sub>, and then  $\gamma$ -Al<sub>2</sub>O<sub>3</sub>, indicating that the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> is the most stable phase as a dense bulk crystal. The surface energies of (100), (110), and (111) facets of the cubic  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> and  $\delta'$ -Al<sub>2</sub>O<sub>3</sub> were calculated, and the (0001), ( $1\bar{1}00$ ), and ( $11\bar{2}0$ ) facets were calculated for hexagonal α-Al<sub>2</sub>O<sub>3</sub> (Supplementary Fig. 20). The surface energy of Al<sub>2</sub>O<sub>3</sub> is affected by the hydroxylated/anhydrous surface states<sup>13,37</sup>. The FT-IR spectra show that the intermediate  $\delta'$ -Al<sub>2</sub>O<sub>3</sub> phase and the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> product are anhydrous (Fig. 2e and Supplementary Fig. 21a), while the γ-Al<sub>2</sub>O<sub>3</sub> is hydroxylated. The TGA shows that the surface hydroxyl group density is ~2 OH nm<sup>-2</sup> (Supplementary Fig. 21b); hence, the surfaces of y-Al<sub>2</sub>O<sub>3</sub> were modeled with this hydroxyl density (Supplementary Note 3 and Supplementary Figs. 22-23). It is found that the surface energy is opposite to the bulk energy, where γ-Al<sub>2</sub>O<sub>3</sub> has the lowest surface energy, followed by  $\delta'$ -Al<sub>2</sub>O<sub>3</sub>, and then  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> (Fig. 4a).

The nanocrystal shapes of the three phases were optimized by the Wulff theorem $^{38}$ , and the specific surface area and total energy were calculated (Supplementary Note 4 and Supplementary Fig. 24). The surface energy difference regulates the thermodynamic stability of the three Al<sub>2</sub>O<sub>3</sub> phases (Fig. 4b), driving the phase transformation from  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> to  $\delta^{\zeta}$ -Al<sub>2</sub>O<sub>3</sub> and then to  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> phase as the surface area decreases, consistent with the experimental observation (Fig. 1b). When smaller than a surface area of -93 m² g⁻¹, or larger than a particle size of -17 nm, the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> phase becomes more stable than the  $\delta^{\zeta}$ -phase. Hence, this particle size is suggested as the thermodynamic limit for the synthesis of anhydrous  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> by a thermal process that involves an intermediate  $\delta^{\zeta}$ -phase. The particle size of  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> (-23 nm)

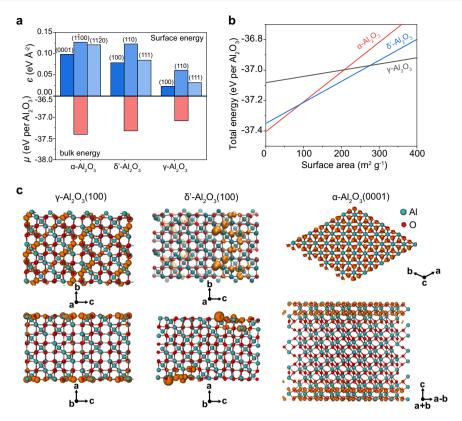


Fig. 4 | Topotactic phase transformation process revealed by density functional theory (DFT) calculations. a The bulk energies ( $\mu$ , eV per Al<sub>2</sub>O<sub>3</sub>) and the surface energies ( $\epsilon$ , eV Å<sup>-2</sup>) of representative crystal surfaces for the three Al<sub>2</sub>O<sub>3</sub> phases. The δ'Al<sub>2</sub>O<sub>3</sub> and α-Al<sub>2</sub>O<sub>3</sub> phases are anhydrous and the γ-Al<sub>2</sub>O<sub>3</sub> phase is hydroxylated with 2 OH nm<sup>-2</sup>. b The total energy (the sum of bulk energy and

surface energy) of the  $Al_2O_3$  nanocrystals of three phases as plotted against the specific surface area.  $\boldsymbol{c}$  The contour plots of partial charge density at the highest bands (0.3 eV below the Fermi levels) of the anhydrous surface states of  $\gamma$ -  $Al_2O_3(100),\,\delta$ -  $Al_2O_3(100),$  and  $\alpha$ -  $Al_2O_3(0001)$  from top view (top) and lateral view (bottom).

synthesized by PDC approaches the thermodynamically limited value, and smaller than those obtained by most other thermal processes (Supplementary Table 7). The ultrafast, pulsed, and low-temperature PDC process to a large extent avoids mass transfer and grain coarsening during the phase transformation process. After adding the entropic contribution, which is minor compared to enthalpy, the free energy vs. surface area is plotted which leads to the same conclusions (Supplementary Note 5 and Supplementary Fig. 25). We note that the energy diagram is merely dependent on the specific surface area and irrelevant to the pore features.

To gain insight into the structural origin of the phase-dependent bulk and surface energy, the partial charge density contour at the highest bands (0.3 eV below the Fermi levels) of the surface states of the three  $Al_2O_3$  phases were plotted (Fig. 4c). All of the surface atoms on  $\alpha\text{-}Al_2O_3(0001)$  are active, while the sites with missing Al atoms on the  $\delta\text{-}Al_2O_3(100)$  and  $\gamma\text{-}Al_2O_3(100)$  surfaces are relatively active (Fig. 4c, top). Closer analysis indicates that the active states go deep into the bulk for the  $\delta\text{-}Al_2O_3(100)$  and  $\gamma\text{-}Al_2O_3(100)$  but not for  $\alpha\text{-}Al_2O_3(0001)$  (Fig. 4c, bottom). This explains the bulk as well as surface energy sequences of the three  $Al_2O_3$  phases, and identifies the Al vacancies in  $\gamma$ - and  $\delta\text{-}$ -phases as the structural origin of their thermodynamic stability/instability vs. the  $\alpha\text{-}$ -phase. The calculation explicitly shows that the phase transformation is thermodynamically driven by the surface energy differences among the three  $Al_2O_3$  phases.

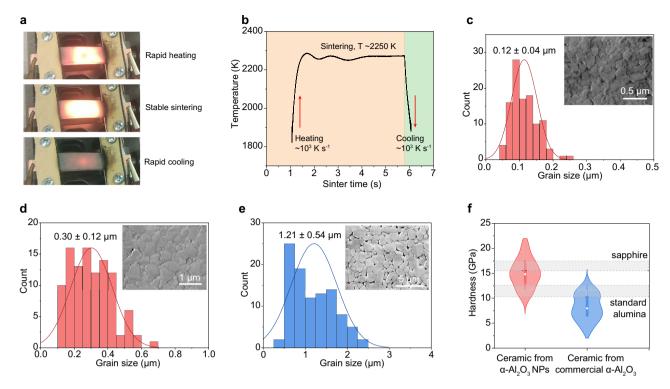
To further verify this, we conducted ab initio dynamic simulations of phase transformation between  $\gamma$ - and  $\alpha$ -phases (Supplementary Note 6 and Supplementary Table 8). Owing to the limitation of current calculation capability, we were unable to depict the full diagram of the phase transformation details. Nevertheless, we found that when the particle size is smaller, the high surface energy of  $\alpha$ -phase drives its

transformation to a structure with typical local order features of  $\gamma\text{-phase}$  (Supplementary Figs. 26 and 27). In contrast, based on the bulk crystal model, the high bulk energy of  $\gamma\text{-phase}$  drives its transformation to a structure with higher coordination numbers of Al and O, approaching those of the  $\alpha\text{-phase}$  (Supplementary Fig. 28 and Supplementary Table 9). Thus, the phase transition from  $\alpha\text{-}$  to  $\gamma\text{-phase}$  is a surface-initiated process while the transition from  $\gamma\text{-}$  to  $\alpha\text{-phase}$  is a bulk defect-initiated process, consistent with the calculated energy landscape (Fig. 4a-b) and the electronic structure analysis (Fig. 4c).

#### Sintering of nanograined alumina ceramics

One prominent application of  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> NPs is as the precursor for sintering high-strength nanograined alumina ceramics. The typical alumina ceramics sintering processes occur under high-pressure and high-temperature (HP-HT) conditions, such as hot isostatic pressing<sup>39</sup>, spark plasma sintering<sup>18</sup>, and pulse electric current sintering<sup>40</sup>. The high pressure, usually several GPa, retains the grain growth and advances densification<sup>41</sup>, which is critical for dense ceramic sintering using coarse grained precursors. The use of additives such as MgO can retard grain growth and advance densification<sup>42,43</sup>. However, the HP-HT process is not suitable for complex structures. The nanocrystalline precursors could undergo the pressureless sintering yet it would suffer from an elevated sintering temperature and prolonged time (>10 h)<sup>6,44,45</sup>. Spark plasma sintering that enables a heating rate up to 600 °C min<sup>-1</sup> has been used for alumina sintering<sup>46</sup>. Very recently, an ultrafast high-temperature sinter method<sup>47</sup> with heating rate up to 10<sup>4</sup> °C min<sup>-1</sup> based on direct current heating is reported for the rapid screening of ceramics.

Here, also based on the Joule heating technique, we developed the alternative current sintering (ACS) process for pressureless,



**Fig. 5** | **Sintering of the nanograined alumina ceramics. a** Picture of the carbon papers during heating, sintering, and cooling. **b** Real-time temperature measurement during the alternating current sintering (ACS) process. **c** Grain size distribution of the alumina ceramic by ACS sintering using the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> NPs as precursor. Inset, scanning electron microscopy (SEM) image of the ceramic. **d** Grain size distribution of the alumina ceramic by two-step pressureless sintering (TS-PS) using the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> NPs as precursor. Inset, SEM image of the ceramic. **e** Grain size

distribution of the alumina ceramic by TS-PS using the commercial  $\alpha\text{-}Al_2O_3$  nanopowders as precursor. Inset, SEM image of the ceramic. **f** Hardness distribution of the alumina ceramics by TS-PS process using the  $\alpha\text{-}Al_2O_3$  NPs (red) and commercial  $\alpha\text{-}Al_2O_3$  nanopowders (blue) as precursors. The dot within the box indicates the median, and the range indicates the 1.5IQR. The hardness of commercial standard alumina  $^{46}$  (10.5–12.7 GPa, lower band) and single-crystal sapphire  $^{50}$  (15.2–17.4 GPa, upper band) are labeled as reference.

ultrafast sintering of alumina ceramics. The ACS system can provide stable and high-energy output with voltages up to 63 V and currents up to 100 A (Supplementary Fig. 29a), making it suitable for the sintering of structural ceramics. Two separated, highly graphitized carbon papers connected to electrodes were used as the heating elements (Supplementary Fig. 29b). The α-Al<sub>2</sub>O<sub>3</sub> NPs, mixed with polyethylene glycol (PEG) binder<sup>48</sup>, were pressed at 500 MPa into green body (See details in Methods). Commercial α-Al<sub>2</sub>O<sub>3</sub> nanopowders (APS ~300 nm) were used as a control. After removal of the binder (5 °C min<sup>-1</sup> to 500 °C for 2 h hold; in air), the green body was put between the carbon papers and under the ACS process at ~15 V (Fig. 5a). The temperature was recorded by fitting the blackbody radiation (Supplementary Fig. 30). The temperature rapidly ramped up to ~2250 K with a heating rate of ~103 K s<sup>-1</sup>. After stable sintering for 5 s, the sample cooled also with a rapid cooling rate of  $\sim 10^3 \,\mathrm{K \, s^{-1}}$  (Fig. 5b).

Even after ACS sintering for only 1 min, the relative density of the ceramics from the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> NPs precursor reaches -97%, higher than that from the commercial  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> nanopowders at -93% (Supplementary Fig. 31 and Supplementary Table 10). The microstructure of the alumina ceramics from  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> NPs by scanning electron microscopy (SEM) showed the average grain size at -0.12  $\mu$ m (Fig. 5c); in comparison, the alumina ceramics sintered from the commercial  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> nanopowders exhibited high residual porosity with grain size of ~1.15  $\mu$ m (Supplementary Fig. 32), demonstrating that the sinter was in its initial stage. These results shows that the ultrafine particle size of the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> NPs facilitates the ultrafast sintering, presumably assisted by the grain growth at high temperature<sup>6</sup>. The mechanical properties were measured on the carefully polished ceramics (see details in Methods). Even with only 1 min ACS sintering, the Young's

modulus and hardness of the ceramics from the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> NPs reach ~86.0 GPa and ~6.2 GPa, respectively (Supplementary Fig. 33), higher than those from the commercial  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> nanopowders (modulus ~40.6 GPa, hardness ~4.2 GPa).

To further improve the mechanical properties, the two-step pressureless sintering (TS-PS) using a high-temperature furnace was applied (see details in Method, Supplementary Fig. 34). Similar to the ACS process, the ceramics sintered by TS-PS from the α-Al<sub>2</sub>O<sub>3</sub> NPs precursor have a higher density (~99%, Supplementary Fig. 31) and much finer grain size (~0.30 µm, Fig. 5d) than those from the commercial α-Al<sub>2</sub>O<sub>3</sub> nanopowder precursors (density ~96% shown in Supplementary Fig. 31, and grain size ~1.21 µm shown in Fig. 5e and Supplementary Fig. 35). The mechanical properties of the ceramics were measured (Fig. 5f and Supplementary Fig. 35). The ceramics sintered from α-Al<sub>2</sub>O<sub>3</sub> NPs precursor demonstrate an average Young's modulus of ~179 GPa (Supplementary Fig. 36a), significantly higher than that from the commercial α-Al<sub>2</sub>O<sub>3</sub> nanopowders (~106 GPa). The median hardness of the ceramics sintered from the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> NPs reaches ~15 GPa (Fig. 5f), which is better than the commercial standard alumina (10.5-12.7 GPa, ref. 49), and comparable to the single-crystal sapphire (15.2–17.4 GPa, ref. 50). We note that the mechanical properties of alumina ceramics are susceptible to various processing parameters<sup>51,52</sup>. For example, the hardness values of 15 -20 GPa were achieved by Krell and coworkers<sup>50,52</sup>, who applied delicate specimens fabrication by hot isostatic pressing, pressure filtration, gel casting, etc. While we do not seek to systematically optimize the sintering process here, the density, grain size, and hardness obtained from the as-synthesized α-Al<sub>2</sub>O<sub>3</sub> NPs are already superior to commercial products and comparable to most state-of-the-art reports (Supplementary Table 11).

#### Discussion

To conclude, an ultrafast PDC processing of v-Al<sub>2</sub>O<sub>3</sub> was developed for the synthesis of fully dehydrated α-Al<sub>2</sub>O<sub>3</sub> NPs (~23 nm) at significantly reduced temperature and duration (~573 K, <1s) than previous thermal processes. Numerical simulations reveal that the phase transformation was enabled by the resistive hotspot effect that induces local heating in the PDC process. Being a highly efficient energy supplies technology, Joule heating has a coefficient of performance of 1.0. The localized heating by resistive hotspots in PDC makes the process more effective because most of the electrothermal energy is directly targeted to the phase transformation. The phase transformation synthesis of α-Al<sub>2</sub>O<sub>3</sub> NPs is realized with a low-energy input of ~4.77 kJ g<sup>-1</sup> or \$0.027 kg<sup>-1</sup> in electrical energy cost, which is at least 20× less energy consumptive than a normal furnace annealing process (Supplementary Note 7). Moreover, the PDC process could be scalable by adjusting sample cross-sectional area and the PDC voltage, as suggested by the theoretical analysis of the key parameters (Supplementary Note 8). We demonstrated the synthesis of α-Al<sub>2</sub>O<sub>3</sub> NPs up to 1.4 g-scale per batch within a similar timescale by using a higher PDC voltage (Supplementary Fig. 37).

The PDC process combined with the resistive hotspot effect greatly reduces the required temperature for reactions that should be originally triggered at a high-energy input, serving as an alternative technique for cost-efficient synthesis. We envision that this would be a universal strategy in electrical heating processes whenever a composite (one conductive phase and one non-conductive phase, or two phases with different conductivities) is used. The ACS process with the ultrafast and energy-efficient features, as demonstrated by the sintering of the alumina ceramic within 1 min, could also be promising in the sintering of functional ceramics, porous ceramics, or for materials screening<sup>47</sup>.

#### Methods

#### **Materials**

Commercial  $\gamma$ -Al $_2O_3$  nanopowders (US Nano, 99.99%, hydrophilic, average particle size of 5 nm, made by high-temperature combustion method) were used as the precursors. Prior to use, the starting material is extensively characterized by XRD (Fig. 1c), Rietveld refinement (Supplementary Fig. 6a), FT-IR (Fig. 2e), SEM (Supplementary Fig. 3), TEM (Supplementary Fig. 4), and BET (Supplementary Fig. 5). The XRD refinement shows that the starting material is composed of ~91 wt%  $\gamma$ -Al $_2O_3$  and ~9 wt%  $\gamma$ -Al $_2O_3$  according to the quantitative analysis. The  $\gamma$ -Al $_2O_3$  could be easily decomposed to  $\gamma$ -Al $_2O_3$  by mild calcination (in air, 700 °C for 1 h; Supplementary Fig. 6b). CB (Carbot, BP-2000) was used as the conductive additive. Commercial  $\alpha$ -Al $_2O_3$  nanopowders (US Nano, 99.9%, 300 nm) were used as comparison for alumina ceramics sintering.

#### PDC system and alumina phase transformation process

The electrical diagram of the PDC system is shown in Supplementary Fig. 1a. The  $\gamma\textsc{-Al}_2O_3$  NPs precursor and CB with specific mass ratios were mixed by planetary ball milling (MSE Supplies, PMV1-0.4 L) for 2 h. The precursors (-150 mg per batch) were loaded into a quartz tube with an inner diameter (I.D.) of 8 mm and outside diameter (O.D.) of 12 mm. Graphite rods were used as the electrodes in both ends of the quartz tube. The use of graphite electrodes prevents the contamination of the product. The tube was then loaded on the reaction stage and connected to the PDC system. The resistance was controlled by the compressive force on the electrodes across the sample. The reaction stage was loaded into a plastic vacuum desiccator chamber under a mild vacuum (-10 mm Hg). A capacitor bank with a total capacitance of 0.624 F was charged by a DC supply that can reach voltages up to 500 V. A relay with programmable ms-level delay time was used to control the discharge time.

A variable-frequency drive (VFD) was used to generate pulsed voltage with a frequency that ranges from 0 to 1000 Hz. In the synthesis, a voltage frequency of f=1000 Hz was used. To prevent overheating, the discharge period (or ON state) was set to 20%, and the rest period (or OFF state) was set to 80% (Supplementary Fig. 1b). The heating time was calculated according to the discharge period. Detailed conditions are listed in Supplementary Table 1. After the PDC Joule heating, the apparatus was cooled to room temperature. CAUTION: There is a risk of electrical shock if improperly operated. Safety guidelines can be found in the Supplemental. After PDC Joule heating synthesis, the mixture of  $Al_2O_3$  NPs and CB residues were calcined at  $Al_3O_3$  NPs and CB residues were calcined at  $Al_3O_3$  NPs product.

#### **Alumina ceramics sintering**

**Green body preparation.** PEG was used as the binder<sup>48</sup>. PEG (M.W. 10000) was dissolved into deionized (DI) water at a concentration of 1 wt%. The  $\alpha\text{-Al}_2O_3$  NPs (-20 mg) and PEG were mixed with the PEG ratio of 3 wt%. After drying at 80 °C for 3 h in air, the  $\alpha\text{-Al}_2O_3$  NPs were pressed into pellets (diameter of 5 mm, thickness of -0.5 mm) using a hydraulic press (500 MPa, dwell time of 10 min, Strongway Benchtop 10-Ton Hydraulic Shop Press). The binder was removed by calcining at 500 °C for 3 h in air at a heating rate of 5 °C min^-1 using a Mafu furnace (NEY, MODELS 6-525). The commercial  $\alpha\text{-Al}_2O_3$  nanopoweders (-300 nm) were used as control.

ACS system and ACS sintering process. The electric diagram of the ACS system is show in Supplementary Fig. 29a. The total capacitance of the capacitor bank is 1.5 F. The system was capable of charging to voltage of 0 to 63 V and the current of 0 to 100 A. Two carbon papers (Toray Carbon Paper 060, FuelCellStore) attached to a glass slide were used as the heating element and sample holder (Supplementary Fig. 29b). The resistance of the carbon paper was ~1  $\Omega$ . The  $\alpha$ -Al $_2$ O $_3$  green bodies were put in between the carbon papers, which were connected to the ACS system. The voltage was set to ~15 V and the sintering time was 1 min.

**Two-step pressureless sintering process.** The two-step pressureless sintering (TS-PS) of alumina ceramics was conducted using a Mafu furnace (Carbolite RHF 1500) with the maximum temperature of 1500 °C. In the first step, the sample temperature ramps to 1425 °C with the heating rate of 5 °C min<sup>-1</sup> and is maintained at 1425 °C for 2 h. In the second step, the sample cools to 1350 °C and is maintained for 5 h. Then, the sample slowly cools to room temperature (Supplementary Fig. 34).

**Density measurement.** The densities of the samples, including the green bodies and the sintered ceramics, are measured by Archimedes' method.

Ceramics surface polishing. Prior to the mechanical property measurement, the sintered ceramic samples were carefully ground and polished using a wafer polisher (MultiPrepTM Precision Polishing System). Diamond lapping film with grit of 15, 9, 6, 3, 1, and 0.1  $\mu m$  were sequentially used. After the polishing process, mirror-like surfaces were obtained for the ceramics sintered from the  $\alpha\text{-Al}_2O_3$  NPs precursors. In contrast, we cannot obtain mirror-like surface for the ceramics sintered from commercial  $\alpha\text{-Al}_2O_3$  nanopowders precursors with the same polishing process.

**Mechanical properties measurement**. The Young's modulus and Vickers hardness of ceramics were measured using a Hysitron TI 980 Tribolndenter. A diamond tip was used as the indenter. 25 curves were measured on each kind of ceramic sample surfaces to account for the deviation caused by residual porosity and heterogeneities. The

reduced Young's moduli were measured, and the Young's moduli were calculated by Eq. (4),

$$\frac{1}{E_{\rm r}} = \frac{1 - v^2}{E} + \frac{1 - v_{\rm i}^2}{E_{\rm i}} \tag{4}$$

where  $E_r$  is the reduce Young's modules of the sample,  $E_i$  is the Young's modules of the sample,  $E_i$  is the Poisson's ratio of the sample ( $E_i = 0.22$ ),  $E_i$  is the Young's modules of the indenter ( $E_i = 1220$  GPa), and  $E_i$  is the Poisson's ratio of the indenter ( $E_i = 0.2$ ).

**Microstructure characterization**. For the ceramics sintered by the TS-PS process, the polished samples were thermally etching at 1400 °C using a furnace (Carbolite RHF 1500) for 30 min for SEM imaging<sup>20</sup>. For the ceramics sintered by the ultrafast ACS process, the thermal etching process is not suitable since it may change the original microstructures; hence, these samples were cracked and the fracture surfaces were characterized by SEM<sup>42</sup>. The grain size distribution is obtained by measuring 100 grains using ImageJ.

#### Characterization

SEM images were obtained using a FEI Helios NanoLab 660 DualBeam SEM system at voltage of 15 kV and beam current of 100 pA. The element maps by EDS were obtained on a FEI Quanta 400 ESEM FEG system with an EDS detector (Oxford Instrument). The Raman spectra were acquired using a Renishaw Raman microscope (laser wavelength of 532 nm, laser power of 5 mW, and lens of 50×). XRD was collected by using a Rigaku Smartlab II system configured with a Cu Kα radiation  $(\lambda = 1.5406 \text{ Å})$ . The Rietveld refinement was conducted using the GSAS-II software<sup>53</sup>. Rwp values <5% were achieved to secure good convergence. XPS analyses were conducted using a PHI Quantera XPS system under a base pressure of  $5 \times 10^{-9}$  Torr. Elemental spectra were collected using a step size of 0.5 eV with the pass energy of 26 eV. All of the XPS spectra were calibrated by using the standard C 1s peak at 284.8 eV. TEM images, selected area electron diffraction (SAED), and NBD patterns were taken on a IEOL 2100 field emission gun transmission electron microscope under the voltage of 200 kV. BET measurements were carried out on a Quantachrome Autosorb-iQ3-MP/Kr BET Surface Analyzer by using N<sub>2</sub> as the adsorption/desorption gas at 77 K. FT-IR spectra were obtained using a Nicolet FT-IR Infrared Microscope. TGA measurement for carbon black removal was conducted in air at a heating rate of 10 °C min<sup>-1</sup> by using a Q-600 Simultaneous TGA/DSC from TA instruments. TGA measurement of the surface hydroxyl coverage on  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> was conducted in N<sub>2</sub> at a heating rate of 10 °C min<sup>-1</sup> by using a Mettler Toledo TGA/DSC 3+ system from METTLER TOLEDO. The calcination in air (700 °C for 1 h) to remove the carbon black for purification of α-Al<sub>2</sub>O<sub>3</sub> NPs was conducted using a Mafu furnace (NEY 6-160A). The DLS measurement was performed using a NanoSight NS300 system. Prior to measurement, the samples were dispersed in DI water and ultrasonicated using a cup horn sonicator for 0.5 h.

#### **Temperature measurement**

For the PDC induce phase transformation process, the temperature was measured using an IR thermometer (Micro-Epsilon) with temperature measurement range of 200 to 1500 °C. The thermometer is connected to LabView software by using a Multifunction I/O (NI USB-6009) for real-time temperature recording (Supplementary Fig. 14a). Prior to use, the temperature was calibrated on the thermometer.

For the ACS of alumina ceramics process, the temperature was measured by fitting the blackbody radiation of the sample using a homemade, time-resolved spectrometer (Supplementary Fig. 30). The light emitted from the sample during Joule heating was collected by a 16-channel photomultiplier tube (PMT) array, with the spectrum range of 640–1000 nm. The sampling rate was 10 kHz, allowing

for the temporal resolution of  $-100\,\mu s$ . The emission spectra were then fitted to the blackbody radiation to obtain the temperature (T) using Eq. (5),

$$B_{\lambda}(\lambda, T) = \gamma \frac{2hc^2}{\lambda^5} \frac{1}{e^{hc/\lambda k_B T} - 1}$$
 (5)

where  $\lambda$  is the wavelength,  $\gamma$  is a constant introduced for fitting, h is the Planck constant, c is the speed of light in vacuum, and  $k_{\rm B}$  is the Boltzmann constant.

#### **DFT** calculations

The DFT method<sup>54</sup> was used as implemented in the Vienna ab initio Simulation Package (VASP)<sup>55</sup>. A plane wave expansion up to 500 eV is employed in combination with an all-electron-like projector augmented wave (PAW) potential<sup>56</sup>. Exchange-correlation is treated within the generalized gradient approximation (GGA) using the functional parameterized by Perdew, Burke, and Ernzerhof<sup>57</sup>.

Bulk energy calculation. We first calculated the three bulk crystals,  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>,  $\delta$ '-Al<sub>2</sub>O<sub>3</sub>, and  $\gamma$ -Al<sub>2</sub>O<sub>3</sub>. The  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> is the ground-state structure with a hexagonal R-3c symmetry. Both  $\delta'$ -Al<sub>2</sub>O<sub>3</sub> and  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> structures are derived from the spinel MgAl<sub>2</sub>O<sub>4</sub> structure with the space group of Fd-3m and a cubic unit cell (a = 8.17Å) of 4-Al<sub>4</sub>O<sub>8</sub> layers and 8 Mg atoms in between. The Al atoms are octahedrally coordinated and Mg atoms are tetrahedrally coordinated. By replacing Mg with Al atoms, the obtained spinel Al<sub>3</sub>O<sub>4</sub> has 32 O atoms, 16 Al atoms at the octahedral sites, and 8 Al atoms at the tetrahedral sites. A  $1 \times 1 \times 3$  supercell (Al<sub>72</sub>O<sub>96</sub>) is constructed of 12-Al<sub>4</sub>O<sub>8</sub> layers in z direction with 24-tetrahedral Al atoms between the layers. Removing 8 out of the total 72 Al atoms in the supercell yields the  $Al_2O_3$  formula<sup>58</sup>. The  $\delta'$ -phase is formed by removing an octahedrally coordinated Al atom from each Al<sub>4</sub>O<sub>8</sub> layers, with one out of every three layers being skipped. The y-phase is formed by removing tetrahedrally coordinated Al atoms in such a way that every two of the three layers are skipped. Periodic boundary conditions are applied to the unit cell in all three dimensions, with the size of the unit cells or lattice constants being optimized. The Brillouin zone integrations are performed using Monkhorst-Pack type meshes<sup>59</sup>, with sufficient meshes of k-points chosen so that the energy and lattice constant are fully converged. All structures are fully relaxed when the maximum force on each atom is smaller than 0.01 eV Å<sup>-1</sup>. The results of calculated bulk properties of the three phases are summarized in Supplementary Table 6.

Surface energy calculation. We calculated the surface energies of both pristine and OH-adsorbed surfaces. For the OH-adsorbed surface, H-adsorption exists to balance the charge<sup>37</sup>. For the hexagonal  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> phase, three surfaces of (0001), (1 $\bar{1}$ 00), and (11 $\bar{2}$ 0) are studied and no surface reconstruction is found to further lower the surface energy (Supplementary Fig. 20a). For both δ'-Al<sub>2</sub>O<sub>3</sub> and γ-Al<sub>2</sub>O<sub>3</sub> phases derived from the original cubic spinel structure, the (001), (110), and (111) surfaces are investigated (Supplementary Figs. 20b, c). These defect-rich surfaces are found to have lower-energy reconstructions with rearrangement of surface atoms. The surface slabs (15-23 Å thick) are constructed so that the two surfaces bounded by vacuum are in inversion symmetry. Low-energy structures of adsorption of OH<sup>-</sup> (on the Al sites) and H<sup>+</sup> (on O sites)<sup>37,60</sup> are searched computationally based on the optimized symmetry and homogenous distribution. The optimized atomic structures of y-Al<sub>2</sub>O<sub>3</sub> surfaces with ~2 OH nm<sup>-2</sup> are shown in Supplementary Fig. 23. The vacuum layers between the slabs are chosen to be 10 Å or thicker.

The surface energy ( $\epsilon$ , in eV Å<sup>-2</sup>) is calculated by Eq. (6),

$$\epsilon = (E_{\text{slab}} - N\mu - hE_{\text{H},O})/(2S), \tag{6}$$

where  $E_{\text{slab}}$  is the total energy of the slab, N is the total number of  $Al_2O_3$  units in the slab,  $\mu$  is the bulk energy in eV per  $Al_2O_3$ , h is the number of water molecules adsorbed in the form of split  $OH^-$  and  $H^+$ ,  $E_{H_2O}$  is the total energy of a water molecule, and S is the surface area of one side of the slab. The calculated surface energies of typical crystal surfaces for the three  $Al_2O_3$  phases, both with and without surface OH group adsorption, are shown in Supplementary Table 6.

Nanocrystal shape optimization based on Wulff construction. Strict analysis of formation energy of nanocrystals is based on the generalized Wulff theorem that considers bulk energy, surface energy, edge energy, and vertex energy<sup>61</sup>. However, direct optimization of the Wulff energy is practically forbidden because in most cases it is not feasible to determine the detailed atomic structures of the reconstructed high-index surfaces, edges, vertices. Therefore, the more practical and typical method involves a limited rational search of hypothetical structures of nanocrystals that potentially have the lowest energy<sup>62</sup>. In this study, we optimize the polyhedral shape of nanoparticles based on the surfaces listed in Supplementary Table 6. Optimization of the shape of nanocrystals is based on the generalized Wulff theorem<sup>61</sup>,

$$E_{\text{total}} = \mu + \sum (S_i \epsilon_i + l_j e_j + v_k) / N$$
 (7)

where  $E_{\rm total}$  is the total energy,  $\mu$  is the bulk energy per formular unit (or atom), N is the total number of formular unit in the particle,  $S_i$  is the area of the  $i^{\rm th}$  facet of the nanocrystal,  $\epsilon_i$  is the surface energy area of the  $i^{\rm th}$  facet of the nanocrystal,  $l_j$  is the length of the  $j^{\rm th}$  edge,  $e_j$  is the edge energy, and  $v_k$  is the energy of the  $k^{\rm th}$  vertex. Then the Eq. (7) can be further simplified by considering only the most probable facets with relatively lower indexes because the higher-index facets normally have much higher energy or complicated reconstruction, which make them unlikely occur in nanocrystals. The detailed shape optimization processes of the three  ${\rm Al_2O_3}$  phases are shown in Supplementary Note 4.

**Plotting total energy vs. specific surface area.** The total energy  $(E_{\text{total}})$  of the nanocrystal is expressed as Eq. (8),

$$E_{\text{total}} = (N\mu + \sum_{i} S_{i} \epsilon_{i})/N$$
 (8)

where N is the total number of  $Al_2O_3$  units,  $\mu$  is the bulk energy in eV per  $Al_2O_3$ ,  $S_i$  is the area of each facet of the nanocrystal,  $\epsilon_i$  is the surface energy in eV Å<sup>-2</sup>, and  $N = V \times n$ , where n is the density, and volume V and  $S_i$  are the functions of the structural parameters of the polyhedrons. The specific surface area is calculated by Eq. (9),

surface area = 
$$\frac{\sum S_i}{N(2A_{AI} + 3A_0)u}$$
 (9)

where  $S_i$  is the area of each facet of the nanocrystal, N is the total number of  $Al_2O_3$  units,  $A_{Al}$  is the atomic mass of Al (26.9815),  $A_O$  is the atomic mass of O (15.9994), and u is the atomic mass (1.6605 ×  $10^{-24}$  g). The total energy and the specific surface area of the three  $Al_2O_3$  phases are calculated based on the optimized nanocrystal shapes, and their detailed expressions are shown in Supplementary Note 4. The energy of the nanocrystals is thus plotted against surface area (Fig. 4b).

#### Data availability

The data supporting the findings of this study are available within the article and its Supplementary Information. The source data generated in this study have been deposited in the Zenodo database under

https://doi.org/10.5281/zenodo.6850605. Other relevant data are available from the corresponding authors upon request. Source data are provided with this paper.

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#### Acknowledgements

We thank Dr. Helge Gonnermann of Rice University for allowing us to use the FEM simulation software, and Dr. Bo Chen of Rice University for helpful discussion of the XPS results. The funding of the research is provided by Air Force Office of Scientific Research (FA9550-19-1-0296, J.M.T.), the Office of Naval Research (N00014-18-1-2182, B.I.Y.), and the U.S. Army Corps of Engineers, ERDC (W912HZ-21-2-0050, B.I.Y., Yufeng

Z., and J.M.T.). The authors acknowledge the use of the Electron Microscopy Center (EMC) at Rice University. The characterization equipment used in this project is, in part, from the Shared Equipment Authority (SEA) at Rice University.

#### **Author contributions**

B.D. and J.M.T. conceived the idea to use the PDC for alumina phase transformation. B.D. conducted the synthesis and characterizations with the help of E.A.M., Z.W., J.C., and R.A.C. B.D. conducted the numerical simulation. D.X.L. designed the PDC and ACS systems and the temperature measurement setups. B.D. conducted the ceramics sintering with the help of C.K. J.Z. and Yuji Z. helped with the ceramic polishing. P.A., B.Z., and J.L. conducted the mechanical measurement. Yufeng Z. and B.I.Y. conducted the DFT calculation. B.D., Yufeng Z., and J.M.T. wrote the manuscript. All aspects of the research were overseen by J.M.T. All authors discussed the results and commented on the manuscript.

#### **Competing interests**

Rice University owns intellectual property on the PDC strategy for the phase transformation synthesis of corundum nanoparticles and ACS sintering process. The authors declare no other competing interests.

#### **Additional information**

**Supplementary information** The online version contains supplementary material available at

https://doi.org/10.1038/s41467-022-32622-4.

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**Peer review information** *Nature Communications* thanks the anonymous reviewers for their contribution to the peer review of this work.

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